Numerical modelling of past retreat and future evolution of Chhota Shigri glacier in Western Indian Himalaya

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ABSTRACT. The history of glacier length fluctuations serves as a reliable indicator of the past climate. In this paper, a numerical flowline model has been used to study the relationship between length variations of Chhota Shigri glacier and local climate since 1876. The simulated front positions of Chhota Shigri glacier are in agreement with those observed. After a successful simulation of the past retreat, the model was also used to predict future evolution of the glacier for the next 100 years under different climatic scenarios. These simulations indicate that the Chhota Shigri glacier may lose ~90% of its present volume by 2100 if the local temperature increases by 2.4 K, and for a temperature rise of 5.5 K, the glacier loses almost all its volume.

Keywords: climate change, glacier flow, glacier mass balance, ice thickness measurements, remote sensing

INTRODUCTION

With an ice coverage of 40.775 km², the Himalaya-Karakoram region has one of the largest glacier coverages outside the Polar regions (Bolch and others, 2012; Frey and others, 2014). Livelihood of a sizeable part of the population is dependent on glacier fed rivers that originate from the Himalaya. Recent investigations suggest that the majority of Himalayan glaciers are retreating due to climate change (Kulkarni and others, 2011; Kulkarni and Karyakarte, 2014). Therefore, in order to assess their sustainability as well as future evolution, it is important to explore the relationship between glacier front fluctuations and climate change. A record of glacier mass balance is an important tool for such studies. However, owing to harsh climatic conditions and difficult terrain in this region, mass-balance investigations have been difficult and also are limited to few selected glaciers (Berthier and others, 2007). Apart from mass balance, topographic maps and glacier retreat records do give some insight into changes in glacier length and surface area, but the state of balance or imbalance with respect to local climate cannot be inferred (Oerlemans, 2001). In such a situation, numerical modelling of glacier flow can be a feasible alternative to in-situ measurements. Numerical flowline modelling has been successfully used for estimating past retreat as well as future evolution of glaciers located in different climatic conditions, around the world (Stroeven, 1996; Oerlemans, 1997; Schmeits and Oerlemans, 1997; De Smedt and Pattyn, 2003). Unfortunately, such flowline studies have not been carried out in the Indian Himalaya.

In this paper, we use a numerical flowline model to simulate the past retreat and future evolution of Chhota Shigri glacier, located in the Indian Himalaya. The glacier was considered for our analyses, mainly because it has been chosen as the benchmark glacier in the western Himalaya; is easily accessible; and most importantly, it is located in an area influenced by Indian monsoon in summer and mid-latitude westerlies in winter (Wagnon and others, 2007). A wide accumulation area and a relatively narrow ablation zone renders the glacier sensitive to climate change. Also, among other Himalayan glaciers Chhota Shigri has one of the longest records of continuous in-situ mass-balance estimates spanning 2002–14 (Azam and others, 2016).

STUDY AREA

The Chhota Shigri glacier is located on the Chandra river basin in Himachal Pradesh (Fig. 1). The valley glacier is oriented north-south in ablation zone, and the accumulation zone is differently oriented. The glacier geometry is complex with three tributaries. The ice flow in the main trunk is mainly due to ice from the tributary (B) and the main trunk (A) (Wagnon and others, 2007). The ice flow in the accumulation zone has contribution from four small cirques namely b1, b2, b3 and b4 (Fig. 1). Continuous mass-balance observations have been carried out on this glacier since 2002 by the glaciological method (Wagnon and others, 2007; Azam and others, 2016). The observed specific mass balance averaged over every altitudinal band of 50 m is shown as a function of altitude in Figure 2. The specific mass-balance ranges from 1.5 to −2.5 m w.e. a−1. Below an altitude of ~4400 m, the specific mass balance increases due to the presence of debris cover (Azam and others, 2016). Further details regarding specific mass balance of Chhota Shigri glacier can be found in Wagnon and others (2007); Azam and others (2016). In addition, point surface velocity and ice thickness observations are available for the year 2004/05 (Tiwari and others, 2014). The snout position has been monitored from 1962 onwards (Ramanathan, 2011). Since 1995, the glacier has been constantly retreating. From 1972 till 2009, the glacier has retreated by 800–850 m (Kulkarni and others, 2007; Shruti, 2008). The other characteristics of the glacier are listed in Table 1.

METHODOLOGY

Model used

The model used in this study is solely based on mass conservation, and it describes ice flow along the central flowline,
assuming that a glacier cross section can be approximated as a trapezium. A detailed description of the model can be inferred from Oerlemans (1997). For a given glacier cross section, mass conservation can be stated as

\[ \frac{\partial H}{\partial t} = - \frac{1}{w_b + 4H} \left( (w_b + 0.5\lambda H)UH \right) + B, \]

where \( H \) is local ice thickness (m), \( w_b \) the width of cross-section’s base (m), \( \lambda \) parametrizes the orientation of the side walls (Fig. 1), \( B \) is the specific mass balance (m w.e. a\(^{-1}\)), \( t \) the time and \( U \) is vertically averaged mean cross-sectional velocity (m a\(^{-1}\)). Since, the model is 1-D and neglects transverse gradients of velocity, vertically averaged and mean cross-sectional velocity are the same. \( U \) can be further expressed as

\[ U = f_k \left( f_1 Y H^2 \left( \frac{\partial h}{\partial x} \right)^2 + f_2 Y H \left( \frac{\partial h}{\partial x} \right) + f_3 Y H^2 \left( \frac{\partial h}{\partial x} \right)^3 \right), \]

where \( Y = \frac{1}{3} \), \( \rho \) is the ice density and is assumed to have a constant value of 900 kg m\(^{-3}\), \( g = 9.8 \) m s\(^{-2}\), \( f_1 \) and \( f_2 \) represent the deformation and sliding flow parameters and are assumed to be \( 1.9 \times 10^{-24} \) Pa\(^{-3}\) s\(^{-1}\) and \( 5.7 \times 10^{-20} \) Pa\(^{-3}\) m\(^2\) s\(^{-1}\), respectively (Budd and others, 1979), \( f_k \) is a correction factor that parametrizes the effects of lateral and longitudinal drag (Nye, 1965; Adhikari and Marshall, 2011).
Table 1. Chhota Shigri glacier mass-balance information

<table>
<thead>
<tr>
<th>Glacier name</th>
<th>Length (2009)</th>
<th>ELA in 2009</th>
<th>Years of mass-balance available</th>
</tr>
</thead>
<tbody>
<tr>
<td>Chhota Shigri</td>
<td>9</td>
<td>4925 m</td>
<td>2002–14</td>
</tr>
</tbody>
</table>

2012). However, lateral and longitudinal stresses seldom have an impact over decadal or century scale retreat as compared with mass balance (Oerlemans, 2001). In fact on repeating the analyses with \( f_k = 1 \), the results were not different from those that were obtained by assuming \( f_k \) to be a constant but not equal to one. This is probably because as compared with other model uncertainties and the rather simplistic treatment of ice flow, the effect of the correction factor is probably insignificant (Oerlemans, 2001); \( h \) is surface elevation (m), expressed as

\[
h = h_0 + H,
\]

(3)

where \( h_0 \) is bed elevation (m). The model equation used in our analyses can be obtained by combining Eqns (1) and (2):

\[
\frac{\partial H}{\partial t} = -\frac{1}{w_b + \lambda H} \left( D \frac{\partial h}{\partial x} \right) + B,
\]

(4)

where \( D \) is defined as diffusivity \((m^2 \text{ a}^{-1})\) and is expressed as:

\[
D = (w_b + 0.5AH)f_k^2 \left( I_1 YH^2 \left( \frac{\partial h}{\partial x} \right)^2 + I_2 YH \left( \frac{\partial h}{\partial x} \right)^2 \right).
\]

Specific mass-balance \( B \) can be further expressed as (Oerlemans, 2001)

\[
B(h, t) = B_{ref}(h) + b(t) + \delta B,
\]

(6)

where \( B_{ref} \) is the reference specific mass-balance profile and is solely a function of \( h \). \( B_{ref} \) is obtained from glaciological records; \( b(t) \) represents the balance perturbations derived from proxies such as temperature, precipitation, tree rings etc. \( \delta B \) is purely a function of time and is applied irrespective of altitude; \( \delta B \) is a constant whose value is chosen such that the model produces equilibrium states that are very close to current glacier stands. \( \delta B \) is also applied irrespective of altitude.

Estimation of bed topography

Local ice thickness \((H)\) is estimated using the relationship between surface velocity, slope and ice thickness (Paterson, 2010; Gantayat and others, 2014)

\[
u_x = u_b + \frac{2A}{n+1} \left( f \rho g H \sin(\alpha) \right)^n,
\]

(7)

where \( u_x \) and \( u_b \) represent the surface and basal velocities, respectively; \( u_b \) is assumed to be 25% of \( u_x \); \( n = 3 \) is Glen’s constant (Glen, 1955), \( f \) is a correction factor that parameters lateral and longitudinal stress and it can be tuned using published estimates of observed ice thickness (Nye, 1965; Adhikari and Marshall, 2011, 2012). For Chhota Shigri glacier, \( f \) was tuned at three cross sections namely 1, 2, 3 (Fig. 1) where observed ice thickness estimates were available (Azam and others, 2012), \( A \) is Creep parameter = \( 3.2 \times 10^{-23}\text{Pa} - 3\text{s}^{-1} \) and \( \sin(\alpha) \) is the mean surface slope of every glacier area located between two altitudinal bands having 100 m difference in their respective elevations.

In this paper, for estimating \( H \) Eqn (7) was first solved along discrete points of manually digitized branchlines (e.g. Linsbauer and others, 2012), and the resulting ice thickness was spatially interpolated assuming zero ice thickness at the glacier margin. For interpolation, ANUDEM (Australian National University DEM) was used (Hutchinson, 1989). While digitizing individual branchlines, the following steps were taken: (a) a lateral spacing of 200–220 m was maintained between adjacent lines; (b) a minimum distance of 100–150 m was maintained between the branchlines and glacier margin; (c) the branchlines from a tributary were made to gradually merge with the branchlines from the main stream. Bed elevation was then estimated by subtracting the resulting ice thickness map from the local DEM.

Lastly, the bed topography was smoothed using a 3-by-3 filter to remove small undulations. Details regarding the algorithm can be found in Farinotti and others (2017).

The initial bed elevation profile \((h_0)\) and initial base width \((w_b)\) were estimated from the bed topography estimated above. For estimating the value of \( f_k \) (see Eqn (2)), values of \( f \) across cross sections 1, 2 and 3 (Fig. 1) were estimated by tuning modelled ice thickness in Eqn (7) with GPR estimates. The value of \( f_k \) was assumed to be the mean of the above estimated values of \( f \).

NUMERICAL IMPLEMENTATION

Equation (4) was solved numerically along the two central flowlines with a horizontal grid spacing of 100 m (Fig. 1). It was observed that two flowlines would be sufficient to parametrize the ice flow in the Chhota Shigri glacier complex (Fig. 1). A forward explicit finite-difference scheme was used for the analyses.

Addition of ice from cirques to main stream

The accumulation zone of the main stream is partly fed by 3 cirques and the accumulation zone of the tributary by 1 cirque (Fig. 1). Details regarding the bed geometry of the cirques were not known and due to their steep slopes, shallow ice approximation will not be applicable. Therefore, it was decided to estimate only the mass flux from the cirques at every time step. To do so the scheme suggested by Schneits and Oerlemans (1997) was used. First, the ice volume of the cirque at time \( t \) is calculated according to

\[
\frac{dv}{dt} = \sum a(h_i)b(h_i) - \frac{v}{t},
\]

(8)

where \( v \) is the volume of the cirque at time \( t \). Here \( a(h_i) \) and \( b(h_i) \) are the area and the mass balance of the elevation interval (100 m in this study) respectively, centred around \( h_i \); \( i \) is the number of the elevation band. The first term on the right represents the mass gain and the second term represents volume lost to the main stream/tributary. \( t^* \) is defined as the characteristic timescale and has been introduced to take into account owing to the fact that mass flux due to ice flow from the cirques to the main stream/tributary may take several
years. \( t^* \) is defined as
\[
    t^* = \frac{L^*}{U^*},
\]
where \( L^* \) and \( U^* \) are defined as the characteristic length scale and velocity, respectively. \( L^* \) is assumed to be half the cirque’s length. \( U^* \) depends on the mean ice thickness and mean surface slope of the basin. Therefore, \( t^* \) varies from cirque to cirque. A list of \( t^* \) for the 4 cirques has been given in Table 2. Therefore, at every time step \( \Delta t \), an amount of ice volume equal to \((v/\Delta t)\Delta t\) is added from the cirques. The ice volume is then equally distributed over a certain number of grid points on the main stream/tributary (for numerical stability). This is then converted to respective ice thickness and is added at the respective gridpoints.

**Addition of ice from tributary to main stream**

The ice flux from the tributary to the main stream is estimated (Schmeits and Oerlemans, 1997) as follows:
\[
    F = U(wH).
\]

At every time step \( \Delta t \), the flux between the last two gridpoints of the tributary was estimated and was assumed to be a good approximation of the ice flux that is added from the tributary to the main stream. Ice flux times \( \Delta t \) gives the volume that is added to the main stream. If the surface elevation of the tributary was more than that of the main stream, the ice volume was added to the main stream. The volume is equally distributed across the nearest three grid points on the main stream (for numerical stability). The distributed volume is then converted into respective ice thickness, which is added to the main stream at the corresponding gridpoints.

**Estimation of \( \Delta t \) (time step)**

For solving Eqn (4) for a glacier with unit width, a critical time step of 0.01 a was estimated using the CFL condition (Smith, 1978):
\[
    \Delta t \leq (\Delta x)^2/4D_a,
\]
where \( \Delta x \) is the horizontal grid spacing and \( D_a \) is the average diffusivity. Finally, for the simulation of retreat of Chhota Shigri glacier (accounting for the variable width), a step of 0.001 a was chosen because lowering the time step further, did not affect the model results.

**DATA USED AND MODEL_INITIALIZATION**

For solving Eqn (4), we need (a) two central flowlines representing ice flow along main trunk and tributary, (b) surface elevation profile along the flowlines, (c) initial bed elevation profile along the flowlines, (d) initial estimate of \( w_0 \) and \( \lambda^* \), (e) reference mass-balance profiles, and (f) proxy data for simulating retreat.

The two central flowlines were manually digitized such that they were orthogonal to the surface elevation contours (Fig. 1). Surface elevation values required for the model were taken from the local ASTER DEM (downloaded freely from earthexplorer.usgs.gov/). To remove spatial bias, it was co-registered with one of the Landsat 8 images and was then smoothened using a 3-by-3 filter. The DEM obtained now was used in further analyses. The surface elevation values required for Eqn (1) were extracted from DEM using ArcMap 9.2. For estimating initial bed topography, ice thickness was estimated from surface velocities generated using 2 Landsat-8 PAN scenes (Table 3). Surface velocities were estimated using sub-pixel correlation (e.g. Leprince and others, 2007; Gantayat and others, 2014). To tune the ice thickness model, we have published estimates of observed ice thickness across three cross profiles. These observations were taken during October 2009 (Azam and others, 2012). While doing so we assumed that over a period of 5–6 years, the ice thickness and surface elevation values will not change drastically (\(<\)RMS error of the DEM). Mass-balance profile (2002–14) for the main stream was inferred from Azam and others (2016). The balance profile of the tributary was derived by correcting the main stream’s mass-balance profile for altitude (Fig. 2). \( B_{ref} \) for both main stream and tributary was estimated by fitting a second-order polynomial through the data points (Figs 2a, b). For simulating retreat, annual mean temperatures and annual and seasonal accumulated precipitation (since, 1876) from the meteorological observatory at Shimla has been taken. The observatory is located at 31.1048° N, 77.1734° E, at an elevation of 2202 m and is maintained by the Indian Meteorological Department (IMD) (Fig. 1).

Using the above data, zero ice volume as the initial condition, and the flow parameters \( f_1 \) and \( f_2 \) specified above, the model (Eqn (4)) was run by setting \( dB = 0 \). \( dB \) was then adjusted in order to simulate a steady-state glacier having the same length as that of Chhota Shigri in 2009 (9 km) (Wagnon and others, 2007). The flow parameters and the bed geometry were adjusted till the RMS difference between simulated and observed surface elevation in 2009 was minimum as well as until the area-altitude distribution of the simulated glacier matched with that derived from the DEM within acceptable limits (\(<0.1 \text{ km}^2 \) in our case) (e.g. Oerlemans, 2001). While adjusting the bed, care was taken to ensure that the geometry of the cross sections where ice thickness was modelled using GPR measurements was not altered (Fig. 1). While adjusting the flow parameters, due to lack of field estimates of \( f_1 \) and \( f_2 \), it was decided to keep the ratio \( f_1/f_2 \) constant at all times. It was observed that along with the above specified modifications, when \( dB \) was assigned a value of 0.03 and 0.06 m w.e. a\(^{-1}\) for A and B, a steady-state glacier with length corresponding to that in 2009 was simulated. Such an approach for tuning flow parameters and bed geometry has been used in the past (e.g. Adhikari and Huybrechts, 2009). Instead of tuning the flow and bed parameters by simulating a steady state, we also tried to tune the bed and flow parameters by minimizing the RMS difference between observed and a simulated dis-equilibrium 2009 surface elevation values. The dis-equilibrium state was simulated by forcing the model with a
constant balance perturbation of 0.08 m w.e. a\(^{-1}\) from AD1000–1875 and b(t) for the period 1876–2009. The results obtained from both kinds of tuning were very similar.

Finally, the adjusted bed along with the adjusted flow parameters was used in subsequent analyses.

**RESULTS**

**Estimation of bed topography**

The bed topography of Chhota Shigri glacier is shown in Figure 3. The bed elevation varies from ~4200 to 4500 m in the ablation zone and ~5000–5900 m in the accumulation zone. The bed elevations along three cross sections were validated with GPR measurements (Figs 1, 4). The mean and standard deviation of the uncertainty is ~17 and ~16 m, respectively. The empirical Cumulative Distribution Function (CDF) of the uncertainty shows that ~74% of observed ice thickness values were within ±25% uncertainty in the corresponding modelled ice thickness profiles. A comparison with the cross-sections modelled using Gantayat and others (2014) also shows that the proposed modification using branchlines gives better results.

**Estimation of response time**

When the model is run with the adjusted parameters, it takes ~320 years to produce a steady-state glacier. The steady-state glacier is then subjected to various balance perturbations after 500 years of model run and the subsequent response of the glacier length is shown in Figure 5.

The length response time is defined as the time period required to reach \((1 \times 10^{-1})\) of the final response after a stepwise perturbation in climate. The response times were 75 and 67 years for balance perturbations of ~0.1 and ~0.2 m w.e. a\(^{-1}\), respectively. For positive perturbations of 0.1, 0.2 m w.e. a\(^{-1}\) the response times were 58 and 67 years, respectively. The response times for positive perturbations is slightly less than that for negative perturbations. This is because during retreat the glacier front enters a bed with a relatively smaller slope as compared with that during advance. Similarly, the volume response times for the same set of balance perturbations were 35, 36 years for positive perturbations (0.1, 0.2 m w.e. a\(^{-1}\)) and 38, 40 years for negative perturbations (−0.1, −0.2 m w.e. a\(^{-1}\)). Therefore, the average volume response time leads the average length response time by ~30 years. The volume response time was also estimated according to Johannesson and others (1989),

\[
\tau_{r_{vh}} = \frac{H}{\left|B_{term}\right|},
\]

where \(\tau_{r_{vh}}\) is the volume response time, \(H^*\) is mean ice thickness and \(B_{term}\) is the specific balance at the terminus. \(\tau_{r_{vh}}\) for Chhota Shigri is ~33 years.

The glacier carries memory of the past climate for a time period equal to one or two times the response time (Oerlemans, 2001). Therefore, the length response time gives an estimate of how far in time one must run the model to correctly simulate the historical retreat.

**Simulation of retreat**

The observed retreat record for Chhota Shigri glacier is rather short (48 years). The large response time of the glacier suggests that the model run should start well before the retreat record. In addition, the mass-balance observations span the period 2002–14, which is also quite small. Therefore, in order to reconstruct the mass-balance history over the past two centuries, we used climate proxies such as temperature and precipitation. We decided to adopt a temperature index like approach (e.g. Zuo and Oerlemans, 1997; Adhikari and Huybrechts, 2009).

For reconstructing the past mass-balance history, annual mean temperature and annual and seasonal accumulated precipitation (1876–2009) from the Meteorological station at Shimla, Himachal Pradesh, were used (Fig. 6a). The major reason behind choosing this station is its long records of temperature and precipitation that span over 134 years, which is also roughly double the average length response time and for a successful retreat simulation, the model run must start from a time that is twice the response time (Oerlemans, 1997; Adhikari and Huybrechts, 2009). The other nearby meteorological stations have much shorter records.

Before carrying out further analysis, we chose to fill in the data gaps. For that we estimated the average of five data points prior to the gap and five points later. We assumed that the average of the two averages is a good estimation
A model inter-comparison was done with the ice thickness profiles estimated according to Gantayat and others (2014). (d) The empirical cumulative probability distribution function (CDF) of the difference in observed and modelled ice thickness at three cross sections.

Fig. 5. Response of the steady-state glacier length to different balance perturbations. The response times for positive perturbations are slightly lesser than that for negative perturbations. Average length response time is ∼67 years.

Fig. 4. (a–c) Validation of modelled ice thickness profiles across three cross sections with GPR derived ice thickness (Azam and others, 2012). A model inter-comparison was done with the ice thickness profiles estimated according to Gantayat and others (2014). (d) The empirical cumulative probability distribution function (CDF) of the difference in observed and modelled ice thickness at three cross sections.

for the void. The annual mean temperature showed more correlation with the annual mean mass balance \( (R = 0.5) \) than that between annual accumulated precipitation and annual mean mass balance \( (R = 0.4) \). Due to the limited years of mass-balance measurements (2002–14), a proper forcing function had to be established for simulating retreat. Assuming a linear relationship between past balance perturbations, and temperature and precipitation (e.g. Adhikari and Huybrechts, 2009), \( b(t) \) is expressed as:

\[
b(t) = \beta(\Delta T + \mu_1) + \theta(\Delta P + \mu_2)
\]

where \( \Delta T, \Delta P \) represent temperature and precipitation anomalies by ratio, respectively. Precipitation anomalies by ratio were obtained by dividing annual/seasonal mean precipitation by their corresponding mean for the period 1876–2009. In our analyses, anomalies were estimated considering 1876–2009 as the base period. \( \beta \) and \( \theta \) are defined as the mass-balance sensitivity to temperature (m w.e./°C) and precipitation sensitivity (m w.e./m⁻¹) respectively. Their values are unique for each glacier (Adhikari and Huybrechts, 2009); \( \mu_1 \) and \( \mu_2 \) serve as tuning parameters. Their values are fixed by trial and error so as to get minimum RMS difference between observed and simulated snout positions.

The maximum difference between the observed ELAs as well as the difference between the corresponding annual mass balance was estimated (from 2002 to 2014) from Azam and others (2016). Assuming a vertical lapse rate of −6.5°C km⁻¹, the value of \( \beta \) was estimated to be −0.7 m w.e./°C. The estimated value of \( \beta \) is well within the range of \((-0.52 ≤ −1)\), which was estimated by Azam and others (2016) using a degree day model. Interestingly, the value of −0.7 m w.e./°C was estimated at the ELA by Azam and others (2016). The lapse rate is likely to vary seasonally in a complex terrain like the Himalaya. But since the model is forced with mass-balance changes averaged over a year, a constant annual mean value for the lapse rate was assumed, which has also been used previously (e.g. Adhikari and Huybrechts, 2009).

The winter accumulated precipitation recorded at base camp shows a maximum difference of 290 mm w.e. for the years 2009/10 and 2012/13 (Azam and others, 2016). For estimating \( \theta \), the change in annual mean mass balance for the corresponding years was estimated, \( \theta \) was estimated to be 0.003 m w.e. mm⁻¹. In order to get minimum RMS error between observed and simulated lengths, the values of \( \mu_1 \) and \( \mu_2 \) were estimated to be 0.08 K and 0.07, respectively.

Our simulation of historic retreat started at the year AD 1000 with zero ice volume as the initial condition and 8B was set to zero for both the main stream as well as the tributary. Due to non-availability of meteorological data prior to 1876, we decided to use the period from 1000 to 1875 as a tuning period. From 1876 onwards, the model was forced with \( b(t) \) as shown in Figure 6b. A comparison of the simulated and observed surface elevations of the main stream shows good correlation with an RMS difference of ∼16 m (Fig. 7). The simulated and observed surface elevation profiles of the tributary also match quite well with an RMS error of ∼15 m (Fig. 7). The simulated glacier length fluctuations also match with that observed with an RMS error of ∼150 m (Fig. 8). The simulated glacier in 2009 is longer than the observed length by ∼100 m.
DISCUSSIONS

Characteristics of the simulated 2009 glacier stand

The simulated driving stress is shown in Figure 9. Except for a small region in the altitude band between ∼4900 and 5000 m, the driving stress is in the range 100–200 KPa, which is typically the range in glacier ice. The very high value of the stress ∼250 KPa in the altitudinal range 4900–5000 m is due to the presence of very steep slopes in that part of the glacier bed.

The simulated mean velocity field is also shown in Figure 9. The mean velocity ranged from ∼180 to 10 m \(a^{-1}\). As ice is lost, the velocity vanishes at the top and glacier’s snout. Very high velocity is again found in that altitudinal band of 4900–5000 m due to very high bed slope. The velocity peaks near the two icefalls at ∼6.5 and ∼8 km from the glacier top. High bed slopes result in less ice thickness, thus increasing the sliding component of the mean velocity (Figs 9, 10). A direct comparison of simulated mean velocity profile and that observed is difficult because the field measurements do not allow a calculation of the mass flux through a glacier cross section averaged over several years, which is what the model simulates (Oerlemans, 1997). However, a comparison of surface velocities for the year 2004/05 (Azam and others, 2012) with mean velocity simulated in 2009 shows that the simulated mean velocity is typically in the range 0.7–1 times the surface velocities everywhere, except in the altitudinal band of 4900–5000 m.

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Fig. 6. (a) Meteorological data from Shimla (1876–2009) (a) Annual mean temperature anomalies, (b) Annual mean precipitation anomalies by ratio, (c) Seasonal accumulated precipitation anomalies by ratio. (b) Balance perturbations derived from (a) annual mean temperature anomalies, (b) annual mean temperature and annual mean precipitation anomalies, (c) Annual mean temperature and seasonal mean precipitation anomalies.

Fig. 7. (a) Comparison of simulated and observed surface elevation of A (main stream). The RMS error is ∼16 m. (b) Comparison of simulated and observed surface elevation of B (tributary). The RMS error is ∼15 m. The adjusted bed profile has been shown as a green line.

Fig. 8. Glacier length fluctuations due to mass-balance changes derived from climate forcing (Fig. 6b). The ‘+’ represent the observed front positions. The constant length prior to 1876 is due to the fact that a constant balance perturbation of 0.08 m.w.e. \(a^{-1}\) was used. The RMS difference between observed and simulated glacier lengths is ∼150 m.
Steep bed slopes indicate presence of icefall in the altitudinal band ∼4900–5000. Here, the model over-estimates the velocity. This is due to the simplistic manner in which $U$ has been formulated in Eqn (2). In reality sliding occurs possibly due to the subglacial lubrication by meltwater that is generally absent in the icefalls. As a result, the effect of higher order stresses on sliding becomes dominant (Adhikari and Marshall, 2013). Another reason could be due that the Chhota Shigri glacier is variously oriented in the accumulation zone (Wagnon and others, 2007), the flowline approximation of the model might not hold true. However, this does not affect the simulation of frontal variation probably because that is mostly governed by dynamics in the ablation area (De Smedt and Pattyn, 2003).

Analysis of past retreat and future projections

The simulated glacier length increased from 1900 to 1950 and remained constant until ∼1991. From 1992 onwards, there has been a continuous retreat; the advance resulted due to a relatively cool period from ∼1900 to 1960 (Figs 6b, 8). The retreat after 1995 seems to be quite rapid when compared with the rate of advance during the period 1900–50. From 1972 to 2006, the modelled retreat was ∼700 m, which is quite close to 850 m, a value that was estimated using satellite imagery (Shruti, 2008). The fast retreat can be attributed to two factors: (a) rise in temperature; (b) presence of steep bed below an altitude of 4200 m. The simulation also suggests that from ∼1972 to 1992, the glacier was in a steady state.

We also tried to determine the dominant climate proxy responsible for retreat. Two experiments were carried out. In the first one, we estimated retreat using only temperature anomalies. In the second experiment, retreat was estimated using both temperature and precipitation anomalies. The results show that precipitation hardly influences the glacier length variations (Fig. 8). Thus, we conclude that historical length variations of Chhota Shigri glacier are mostly influenced by temperature and not precipitation.

After having successfully simulated the historic retreat, four different climatic scenarios under constant precipitation were analysed for predicting the glacier’s future evolution. These climatic scenarios have been named as RCP 2.6, RCP 4.5, RCP 6 and RCP 8.6. They predict a temperature rise of 2.36, 3.49, 3.68 and 5.51°C with respect to 1860–1990 mean, respectively, in the Western Himalaya Karakoram over the next 100 years. Details regarding the scenarios can be inferred from Chaturvedi and others (2014).

The trend of the temperature rise was estimated and was used for model runs after 2009. The results are shown in Figure 11. If the mean temperature for the period 1990–2009 was assumed to be the mean temperature after 2009,
the glacier can be expected to retreat by ~4 km and lose 73% of 2009 volume by AD 2109. Under RCP 2.6 and 4.5 situations, the glacier is expected to retreat by ~6 and 6.6 km as well as lose 92% and 97% of its 2009 volume, respectively. Under RCP 6 conditions, the glacier is expected to retreat by ~7 km. Under RCP 8.5 conditions, a steep decrease in the glacier length is to be expected after ~2090 because the glacier breaks to fragments.

CONCLUSIONS
In this study, we present a numerical ice flow model that was used to simulate the historical retreat and future evolution of Chhota Shigri glacier. The simulation was successful with an RMS error of 150 m between simulated and observed snout positions. This study shows that flowline models can be used to understand the relationship between Himalayan glaciers and climate change. The simulations show that within the next 100 years, the glacier is expected to retreat further by between ~6 and 6.6 km for a 2.36 and 3.49 K rise in temperature, respectively. If the annual mean temperature after 60 years increases by between 2.36 and 3.49 K, the glacier is expected to retreat further by ~6 km. Under RCP 2.6 and 4.5 situations, the glacier is expected to retreat by ~7 km for a 2.36 and 3.49 K rise in temperature, respectively.

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